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GAKKOTOSHO TOKYO JAPAN

ASYMPTOTIC ANALYSIS OF AN ε-STOKES PROBLEM CONNECTING STOKES AND PRESSURE-POISSON PROBLEMS

Kazunori Matsui*

Division of Mathematical and Physical Sciences, Graduate School of Natural Science and Technology, Kanazawa University, Kakuma, Kanazawa 920-1192 Japan

(E-mail: first-lucky@stu.kanazawa-u.ac.jp)

and

Adrian Muntean

Department of Mathematics and Computer Science, Karlstad University, Universitetsgatan 2, 651 88 Karlstad Sweden

(E-mail: adrian.muntean@kau.se)

Abstract. In this Note, we prepare an ε -Stokes problem connecting the Stokes problem and the corresponding pressure-Poisson equation using one parameter $\varepsilon > 0$. We prove that the solution to the ε -Stokes problem, convergences as ε tends to 0 or ∞ to the Stokes and pressure-Poisson problem, respectively.

*Corresponding author

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1 Introduction

Let Ω be a bounded domain in $\mathbb{R}^n (n \geq 2, n \in \mathbb{N})$ with Lipschitz continuous boundary Γ and let $F \in L^2(\Omega)^n, u_b \in H^{1/2}(\Gamma)^n$ satisfy $\int_{\Gamma} u_b \cdot \nu = 0$, where ν is the unit outward normal vector for Γ . The weak form of the Stokes problem is: Find $u_S \in H^1(\Omega)^n$ and $p_S \in L^2(\Omega)/\mathbb{R}$ satisfying

$$\begin{cases} -\Delta u_S + \nabla p_S = F & \text{in } H^{-1}(\Omega)^n, \\ \operatorname{div} u_S = 0 & \operatorname{in} L^2(\Omega), \\ u_S = u_b & \text{on } H^{1/2}(\Gamma)^n. \end{cases}$$
(S)

We refer to [20] for details on the Stokes problem, (i.e. more physical background and corresponding mathematical analysis). Taking the divergence of the first equation, we are led to

$$\operatorname{div} F = \operatorname{div}(-\Delta u_S + \nabla p_S) = -\Delta(\operatorname{div} u_S) + \Delta p_S = \Delta p_S$$
(1.1)

in distributions sense. This is often called pressure-Poisson equation and is used in MAC, SMAC or projection method (cf. [1, 4, 7, 12, 13, 15, 17, 19], e.g.). Bearing this in mind, we consider a similar problem: Find $u_{PP} \in H^1(\Omega)^n$ and $p_{PP} \in H^1(\Omega)$ satisfying

$$\begin{cases} -\Delta u_{PP} + \nabla p_{PP} = F & \text{in } H^{-1}(\Omega)^n, \\ -\Delta p_{PP} = -\operatorname{div} F & \text{in } H^{-1}(\Omega), \\ u_{PP} = u_b & \text{on } H^{1/2}(\Gamma)^n, \\ p_{PP} = p_b & \text{on } H^{1/2}(\Gamma). \end{cases}$$
(PP)

with $p_b \in H^{1/2}(\Gamma)$. Let this problem be called pressure-Poisson problem. This idea using (1.1) instead of div $u_S = 0$ is useful to calculate the pressure numerically in the Navier-Stokes equation. For example, the idea is used in both the MAC, SMAC and projection methods [1, 4, 7, 12, 13, 15, 17, 19]. Dirichlet boundary condition for pressure can be found in many circumstances such as outflow boundary [3, 21]. (See also [5, 6, 16].)

In this Note, we prepare on an "interpolation" between these problems (S) and (PP), i.e. we introduce an intermediate problem: For $\varepsilon > 0$, find $u_{\varepsilon} \in H^1(\Omega)^n$ and $p_{\varepsilon} \in H^1(\Omega)$ which satisfy

$$\begin{pmatrix}
-\Delta u_{\varepsilon} + \nabla p_{\varepsilon} = F & \text{in } H^{-1}(\Omega)^{n}, \\
-\varepsilon \Delta p_{\varepsilon} + \operatorname{div} u_{\varepsilon} = -\varepsilon \operatorname{div} F & \text{in } H^{-1}(\Omega), \\
u_{\varepsilon} = u_{b} & \text{on } H^{1/2}(\Gamma)^{n}, \\
p_{\varepsilon} = p_{b} & \text{on } H^{1/2}(\Gamma).
\end{cases}$$
(ES)

Let this problem be called ε -Stokes problem. In [8, 11, 14], they treat this problem as approximation of the Stokes problem to avoid numerical instabilities. The ε -Stokes problem (ES) formally approximates the Stokes problem (S) as $\varepsilon \to 0$ and the pressure-Poisson problem (PP) as $\varepsilon \to \infty$ (Figure 1). We show here that (ES) is a natural link between (S) and (PP) in Proposition 2.7. The aim of this Note is to give a precise asymptotic estimates for (ES) when ε tends to zero or ∞ .



Figure 1: Sketch of the connections between the problems (S), (PP) and (ES).

2 Well-posedness

2.1 Notation

We set

$$C_0^{\infty}(\Omega)^n := \left\{ f \in C^{\infty}(\Omega)^n \mid \operatorname{supp}(f) \text{ is compact subset in } \Omega \right\}$$
$$L^2(\Omega)/\mathbb{R} := \left\{ u \in L^2(\Omega) \mid \int_{\Omega} u = 0 \right\}.$$

For m = 1 or n, $H^{-1}(\Omega)^m = (H^1_0(\Omega)^m)^*$ is equipped with the norm $||f||_{H^{-1}(\Omega)^m} := \sup_{\varphi \in S_m} \langle f, \varphi \rangle$ for $f \in H^{-1}(\Omega)^m$, where $S_m = \{\varphi \in H^1_0(\Omega)^m \mid ||\nabla \varphi||_{L^2(\Omega)^{n \times m}} = 1\}$. We define $[p] := p - (1/|\Omega|) \int_{\Omega} p$ and $||p||_{L^2(\Omega)/\mathbb{R}} := \inf_{a \in \mathbb{R}} ||p - a||_{L^2(\Omega)} = ||[p]||_{L^2(\Omega)}$ for all $p \in L^2(\Omega)$, where $|\Omega|$ is the volume of Ω .

Let $\gamma_0 \in B(H^1(\Omega), H^{1/2}(\Gamma))$ be the standard trace operator. It is known that (see e.g. [20, pp.10–11, Lemma 1.3]) there exists a linear continuous operator $\gamma_{\nu} : H^1(\Omega)^n \to H^{-1/2}(\Gamma)$ such that $\gamma_{\nu}u = u \cdot \nu|_{\Gamma}$ for all $u \in C^{\infty}(\overline{\Omega})^n$, where ν is the unit outward normal for Γ and $H^{-1/2}(\Gamma) := H^{1/2}(\Gamma)^*$. Then, the following generalized Gauss divergence formula holds:

$$\int_{\Omega} u \cdot \nabla \omega + \int_{\Omega} (\operatorname{div} u) \omega = \langle \gamma_{\nu} u, \gamma_{0} \omega \rangle \quad \text{for all } u \in H^{1}(\Omega)^{n}, \omega \in H^{1}(\Omega).$$

We recall the following Theorem 2.1 that plays an important role in the proof of the existence of pressure solution of Stokes problem; see [18, pp.187–190, Lemme 7.1, l = 0] and [9, pp.111–115, Theorem 3.2 and Remark 3.1 (Ω is C^1 class)] for the proof.

Theorem 2.1. There exists a constant c > 0 such that

$$||f||_{L^{2}(\Omega)} \leq c(||f||_{H^{-1}(\Omega)} + ||\nabla f||_{H^{-1}(\Omega)})$$

for all $f \in L^2(\Omega)$.

The following result follows from Theorem 2.1.

Theorem 2.2. [10, pp.20–21] There exists a constant c > 0 such that

$$||f||_{L^2(\Omega)/\mathbb{R}} \le c ||\nabla f||_{H^{-1}(\Omega)^n}$$

for all $f \in L^2(\Omega)$.

2.2 Well-posedness

Theorem 2.3. For $F \in L^2(\Omega)^n$ and $u_b \in H^{1/2}(\Gamma)^n$, there exists a unique pair of functions $(u_S, p_S) \in H^1(\Omega)^n \times (L^2(\Omega)/\mathbb{R})$ satisfying (S).

See [20, pp.31–32, Theorem 2.4 and Remark 2.5] for the proof.

Theorem 2.4. For $F \in L^2(\Omega)^n$, $u_b \in H^{1/2}(\Gamma)^n$ and $p_b \in H^{1/2}(\Gamma)$, there exists a unique pair of functions $(u_{PP}, p_{PP}) \in H^1(\Omega)^n \times H^1(\Omega)$ satisfying (PP).

Proof. From the second and fourth equations of (PP), $p_{PP} \in H^1(\Omega)$ is uniquely determined. Then $u_{PP} \in H^1(\Omega)^n$ is also uniquely determined from the first and third equations.

Corollary 2.5. If the solution $(u_{PP}, p_{PP}) \in H^1(\Omega)^n \times H^1(\Omega)$ of (PP) satisfies div $u_{PP} = 0$, by Theorem 2.3, $u_S = u_{PP}$ and $p_S = [p_{PP}]$ hold.

Theorem 2.6. For $\varepsilon > 0, F \in L^2(\Omega)^n, u_b \in H^{1/2}(\Gamma)^n$ and $p_b \in H^{1/2}(\Gamma)$, there exists a unique pair of functions $(u_{\varepsilon}, p_{\varepsilon}) \in H^1(\Omega)^n \times H^1(\Omega)$ satisfying the problem (ES).

Proof. We pick $u_1 \in H^1(\Omega)^n$ and $p_0 \in H^1(\Omega)$ with $\gamma_0 u_1 = u_b, \gamma_0 p_0 = p_b$. Since div : $H_0^1(\Omega)^n \to L^2(\Omega)/\mathbb{R}$ is surjective [10, p.24, Corollary 2.4, 2°)] and [20, p.32, Lemma 2.4, Chapter 1], there exists $u_2 \in H_0^1(\Omega)^n$ such that div $u_2 = \operatorname{div} u_1$. We put $u_0 := u_1 - u_2$, and then $\gamma_0 u_0 = u_b$ and div $u_0 = 0$ in Ω . To simplify the notation, we set $u := u_\varepsilon - u_0(\in H_0^1(\Omega)^n), p := p_\varepsilon - p_0(\in H_0^1(\Omega)), f \in H^{-1}(\Omega)^n$ and $g \in H^{-1}(\Omega)$ such that $\langle f, v \rangle = \int_{\Omega} Fv - \int_{\Omega} \nabla u_0 : \nabla v - \int_{\Omega} (\nabla p_0) \cdot v \ (v \in H_0^1(\Omega)^n), \langle g, q \rangle = \int_{\Omega} F \cdot \nabla q - \int_{\Omega} \nabla p_0 \cdot \nabla q \ (q \in H_0^1(\Omega))$. Then we have

$$\begin{cases} \int_{\Omega} \nabla u : \nabla \varphi + \int_{\Omega} (\nabla p) \cdot \varphi = \langle f, \varphi \rangle & \text{for all } \varphi \in H_0^1(\Omega)^n, \\ \varepsilon \int_{\Omega} \nabla p \cdot \nabla \psi + \int_{\Omega} (\operatorname{div} u) \psi = \varepsilon \langle g, \psi \rangle & \text{for all } \psi \in H_0^1(\Omega). \end{cases}$$
(2.2)

Adding the equations in (2.2), we get

$$\begin{pmatrix} u \\ p \end{pmatrix}, \begin{pmatrix} \varphi \\ \psi \end{pmatrix} \rangle_{\varepsilon} = \langle f, \varphi \rangle + \varepsilon \langle g, \psi \rangle.$$

Here, we denote

$$\left(\begin{pmatrix} u \\ p \end{pmatrix}, \begin{pmatrix} \varphi \\ \psi \end{pmatrix} \right)_{\varepsilon} := \int_{\Omega} \nabla u : \nabla \varphi + \varepsilon \int_{\Omega} \nabla p \cdot \nabla \psi + \int_{\Omega} (\nabla p) \cdot \varphi + \int_{\Omega} (\operatorname{div} u) \psi.$$

We check that $(*,*)_{\varepsilon}$ is a continuous coercive bilinear form on $H_0^1(\Omega)^n \times H_0^1(\Omega)$. The bilinearity and continuity of $(*,*)_{\varepsilon}$ are obvious. The coercivity of $(*,*)_{\varepsilon}$ is obtained in the following way: Let ${}^t(u,p) \in H_0^1(\Omega)^n \times H_0^1(\Omega)$. We have the following sequence of inequalities;

$$\begin{pmatrix} u \\ p \end{pmatrix}, \begin{pmatrix} u \\ p \end{pmatrix} \rangle_{\varepsilon} = \int_{\Omega} \nabla u : \nabla u + \varepsilon \int_{\Omega} \nabla p \cdot \nabla p + \int_{\Omega} \operatorname{div}(up)$$

$$= ||\nabla u||_{L^{2}(\Omega)}^{2} + \varepsilon ||\nabla p||_{L^{2}(\Omega)}^{2}$$

$$\geq \min\{1, \varepsilon\}(||\nabla u||_{L^{2}(\Omega)}^{2} + ||\nabla p||_{L^{2}(\Omega)}^{2})$$

$$\geq c \min\{1, \varepsilon\}(||u||_{H^{1}(\Omega)}^{2} + ||p||_{H^{1}(\Omega)}^{2})$$

From now on, let the solutions of (S), (PP) and (ES) be denoted by (u_S, p_S) , (u_{PP}, p_{PP}) and $(u_{\varepsilon}, p_{\varepsilon})$, respectively.

Proposition 2.7. Suppose that $p_S \in H^1(\Omega)$. Then there exists a constant c > 0 independent of ε such that

$$||u_{S} - u_{PP}||_{H^{1}(\Omega)^{n}} \le c||\gamma_{0}p_{S} - p_{b}||_{H^{1/2}(\Gamma)}, \quad ||u_{S} - u_{\varepsilon}||_{H^{1}(\Omega)^{n}} \le c||\gamma_{0}p_{S} - p_{b}||_{H^{1/2}(\Gamma)}$$

In particular, if $\gamma_0 p_S = p_b$, then $p_{PP} = p_{\varepsilon} = p_S$ hold for all $\varepsilon > 0$.

Proof. From (S) and (PP), we have

$$\begin{cases} \int_{\Omega} \nabla(u_S - u_{PP}) : \nabla \varphi = -\int_{\Omega} (\nabla(p_S - p_{PP})) \cdot \varphi & \text{for all } \varphi \in H_0^1(\Omega)^n, \\ \int_{\Omega} \nabla(p_S - p_{PP}) \cdot \nabla \psi = 0 & \text{for all } \psi \in H_0^1(\Omega). \end{cases}$$
(2.3)

Putting $\varphi := u_S - u_{PP} \in H^1_0(\Omega)^n$ in (2.3), we get

$$\begin{aligned} ||\nabla(u_S - u_{PP})||^2_{L^2(\Omega)^{n \times n}} &= -\int_{\Omega} (\nabla(p_S - p_{PP})) \cdot (u_S - u_{PP}) \\ &\leq ||\nabla(p_S - p_{PP})||_{L^2(\Omega)^n} ||u_S - u_{PP}||_{L^2(\Omega)^n}, \end{aligned}$$

and then

$$||u_S - u_{PP}||_{H^1(\Omega)^n} \le c_1 ||\nabla(p_S - p_{PP})||_{L^2(\Omega)^n}$$
(2.4)

follows. We pick up $p_0 \in H^1(\Omega)$ such that $\gamma_0 p_0 = p_b$. From the fourth equation of (PP) and the second equation of (2.3), we obtain $p_{PP} - p_0 \in H^1_0(\Omega)$ and

$$\int_{\Omega} \nabla (p_{PP} - p_0) \cdot \nabla \psi = \int_{\Omega} \nabla (p_S - p_0) \cdot \nabla \psi,$$

and, by Stampacchia Theorem [2, Theorem 5.6], it follows that

$$\min_{\substack{\psi \in H_0^1(\Omega)^n}} \left(\frac{1}{2} \| \nabla \psi \|_{L^2(\Omega)^n}^2 - \int_{\Omega} \nabla (p_S - p_0) \cdot \nabla \psi \right) \\
= \frac{1}{2} \| \nabla (p_{PP} - p_0) \|_{L^2(\Omega)^n}^2 - \int_{\Omega} \nabla (p_S - p_0) \cdot \nabla (p_{PP} - p_0) \\
= \frac{1}{2} \| \nabla p_{PP} \|_{L^2(\Omega)^n}^2 - \frac{1}{2} \| \nabla p_0 \|_{L^2(\Omega)^n}^2 - \int_{\Omega} \nabla p_S \cdot \nabla p_{PP} + \int_{\Omega} \nabla p_S \cdot \nabla p_0.$$

Hence,

$$\frac{1}{2} \|\nabla (p_S - p_{PP})\|_{L^2(\Omega)^n}^2 = \min_{\psi \in H_0^1(\Omega)^n} \left(\frac{1}{2} \|\nabla (p_S - p_0 - \psi)\|_{L^2(\Omega)^n}^2\right).$$

Since γ_0 is surjective and the space $\operatorname{Ker}(\gamma_0) = H_0^1(\Omega), H^1(\Omega)/H_0^1(\Omega)$ and $H^{1/2}(\Gamma)$ are isomorphic, there exists a constant $c_2 > 0$ such that $\|q\|_{H^1(\Omega)/H_0^1(\Omega)} \leq c_2 \|\gamma_0 q\|_{H^{1/2}(\Gamma)}$ for all $q \in H^1(\Omega)$. Hence, we obtain

$$\begin{aligned} \|\nabla(p_S - p_{PP})\|_{L^2(\Omega)^n} &\leq \min_{\psi \in H_0^1(\Omega)^n} \|\nabla(p_S - p_0 - \psi)\|_{L^2(\Omega)^n} \\ &\leq \min_{\psi \in H_0^1(\Omega)^n} \|p_S - p_0 - \psi\|_{H^1(\Omega)} \\ &= \|p_S - p_{PP}\|_{H^1(\Omega)/H_0^1(\Omega)} \\ &\leq c_2 \|\gamma_0 p_S - \gamma_0 p_0\|_{H^{1/2}(\Gamma)}. \end{aligned}$$

Together with (2.4) and the assumption $\gamma_0 p_0 = p_b$, we obtain $||u_S - u_{PP}||_{H^1(\Omega)^n} \leq c_1 c_2 ||\gamma_0 p_S - p_b||_{H^{1/2}(\Gamma)}$.

Let $w_{\varepsilon} := u_S - u_{\varepsilon} \in H_0^1(\Omega)^n$, $r_{\varepsilon} := p_{PP} - p_{\varepsilon} \in H_0^1(\Omega)$. By (S), (PP) and (ES), we have

$$\begin{cases} \int_{\Omega} \nabla w_{\varepsilon} : \nabla \varphi + \int_{\Omega} (\nabla r_{\varepsilon}) \cdot \varphi = -\int_{\Omega} (\nabla (p_S - p_{PP})) \cdot \varphi & \text{for all } \varphi \in H_0^1(\Omega)^n, \\ \varepsilon \int_{\Omega} \nabla r_{\varepsilon} \cdot \nabla \psi + \int_{\Omega} (\operatorname{div} w_{\varepsilon}) \psi = 0 & \text{for all } \psi \in H_0^1(\Omega). \end{cases}$$
(2.5)

Putting $\varphi := w_{\varepsilon}$ and $\psi := r_{\varepsilon}$ and adding two equations of (2.5), we get

$$||\nabla w_{\varepsilon}||^{2}_{L^{2}(\Omega)^{n \times n}} + \varepsilon||\nabla r_{\varepsilon}||^{2}_{L^{2}(\Omega)^{n}} = -\int_{\Omega} (\nabla (p_{S} - p_{PP})) \cdot w_{\varepsilon} \le ||\nabla (p_{S} - p_{PP})||_{L^{2}(\Omega)^{n}} ||w_{\varepsilon}||_{L^{2}(\Omega)^{n}}$$

from $\int_{\Omega} (\nabla r_{\varepsilon}) \cdot w_{\varepsilon} = -\int_{\Omega} (\operatorname{div} w_{\varepsilon}) r_{\varepsilon}$. Thus it leads $||w_{\varepsilon}||_{H^{1}(\Omega)^{n}} \leq c_{3} ||\nabla (p_{S} - p_{PP})||_{L^{2}(\Omega)^{n}}$. Hence we obtain $||u_{S} - u_{\varepsilon}||_{H^{1}(\Omega)^{n}} = ||w_{\varepsilon}||_{H^{1}(\Omega)^{n}} \leq c_{2}c_{3}||\gamma_{0}p_{S} - p_{b}||_{H^{1/2}(\Gamma)}$.

Proposition 2.8. Under the hypotheses of Proposition 2.7, if $\tilde{p} \in H^1(\Omega)$ satisfies $\gamma_0 \tilde{p} = p_b$, then we have

$$||\nabla(\tilde{p} - p_{PP})||_{L^2(\Omega)^n} \le ||\nabla(\tilde{p} - p_S)||_{L^2(\Omega)^n}.$$

Proof. By the second equation of (2.3) and $\tilde{p} - p_{PP} \in H^1_0(\Omega)$, it yields

$$\int_{\Omega} \nabla (p_S - p_{PP}) \cdot \nabla (\tilde{p} - p_{PP}) = 0$$

Hence we obtain

$$\begin{aligned} ||\nabla(\tilde{p} - p_{PP})||_{L^{2}(\Omega)^{n}}^{2} &= \int_{\Omega} \nabla(\tilde{p} - p_{S} + p_{S} - p_{PP}) \cdot \nabla(\tilde{p} - p_{PP}) \\ &\leq ||\nabla(\tilde{p} - p_{S})||_{L^{2}(\Omega)^{n}} ||\nabla(\tilde{p} - p_{PP})||_{L^{2}(\Omega)^{n}}. \end{aligned}$$

Therefore, $||\nabla(\tilde{p} - p_{PP})||_{L^2(\Omega)^n} \le ||\nabla(\tilde{p} - p_S)||_{L^2(\Omega)^n}$ holds.

Remark 2.9. If $p_S \in H^1(\Omega)$, then we have

$$||\nabla (p_{\varepsilon} - p_{PP})||_{L^{2}(\Omega)^{n}} \leq ||\nabla (p_{\varepsilon} - p_{S})||_{L^{2}(\Omega)^{n}}$$

for all $\varepsilon > 0$, (from Proposition 2.8). Hence, if $(\nabla p_{\varepsilon})_{\varepsilon>0}$ converges strongly to ∇p_S in $L^2(\Omega)^n$, then there exists a constant $c \in \mathbb{R}$ such that $u_{PP} = u_S$ and $p_{PP} = p_S + c$, which imply $\gamma_0 p_S = p_b + c$ for some $c \in \mathbb{R}$. In other words, if $p_S \in H^1(\Omega)$ satisfies $\gamma_0 p_S \neq p_b + c$ for all $c \in \mathbb{R}$, then ∇p_{ε} does not converge to ∇p_S in $L^2(\Omega)^n$ as $\varepsilon \to 0$.

3 Links between (ES) and (PP)

Theorem 3.1. There exists a constant c > 0 independent of ε satisfying

$$||u_{\varepsilon} - u_{PP}||_{H^{1}(\Omega)^{n}} \leq \frac{c}{\varepsilon} ||\operatorname{div} u_{PP}||_{H^{-1}(\Omega)}, \quad ||p_{\varepsilon} - p_{PP}||_{H^{1}(\Omega)} \leq \frac{c}{\varepsilon} ||\operatorname{div} u_{PP}||_{H^{-1}(\Omega)}.$$

for all $\varepsilon > 0$. In particular, we have

$$||u_{\varepsilon} - u_{PP}||_{H^{1}(\Omega)^{n}} \to 0, \ ||p_{\varepsilon} - p_{PP}||_{H^{1}(\Omega)} \to 0 \quad as \ \varepsilon \to \infty$$

Proof. From (PP) and (ES), we have

$$\begin{cases} \int_{\Omega} \nabla(u_{\varepsilon} - u_{PP}) : \nabla\varphi + \int_{\Omega} (\nabla(p_{\varepsilon} - p_{PP})) \cdot \varphi = 0 & \text{for all } \varphi \in H_0^1(\Omega)^n, \\ \varepsilon \int_{\Omega} \nabla(p_{\varepsilon} - p_{PP}) \cdot \nabla\psi + \int_{\Omega} (\operatorname{div}(u_{\varepsilon} - u_{PP}))\psi = -\int_{\Omega} (\operatorname{div} u_{PP})\psi & \text{for all } \psi \in H_0^1(\Omega). \end{cases}$$
(3.6)

Putting $\varphi := u_{\varepsilon} - u_{PP} \in H_0^1(\Omega)^n$ and $\psi := p_{\varepsilon} - p_{PP} \in H_0^1(\Omega)$ and adding two equations of (3.6), we obtain

$$||\nabla(u_{\varepsilon} - u_{PP})||^{2}_{L^{2}(\Omega)^{n \times n}} + \varepsilon||\nabla(p_{\varepsilon} - p_{PP})||^{2}_{L^{2}(\Omega)^{n}} \leq ||\operatorname{div} u_{PP}||_{H^{-1}(\Omega)}||\nabla(p_{\varepsilon} - p_{PP})||_{L^{2}(\Omega)^{n}},$$

where we have used $\int_{\Omega} (\nabla (p_{\varepsilon} - p_{PP})) \cdot (u_{\varepsilon} - u_{PP}) = -\int_{\Omega} (\operatorname{div}(u_{\varepsilon} - u_{PP}))(p_{\varepsilon} - p_{PP})$. Thus

$$||\nabla(p_{\varepsilon} - p_{PP})||_{L^{2}(\Omega)^{n}} \leq \frac{1}{\varepsilon} ||\operatorname{div} u_{PP}||_{H^{-1}(\Omega)}$$

follows. In addition, by (3.6) and the Poincaré inequality, we have

$$\begin{aligned} ||\nabla(u_{\varepsilon} - u_{PP})||^{2}_{L^{2}(\Omega)^{n}} &= -\int_{\Omega} (\nabla(p_{\varepsilon} - p_{PP})) \cdot (u_{\varepsilon} - u_{PP}) \\ &\leq ||\nabla(p_{\varepsilon} - p_{PP})||_{L^{2}(\Omega)^{n}} ||u_{\varepsilon} - u_{PP}||_{L^{2}(\Omega)^{n}} \\ &\leq c ||\nabla(p_{\varepsilon} - p_{PP})||_{L^{2}(\Omega)^{n}} ||\nabla(u_{\varepsilon} - u_{PP})||_{L^{2}(\Omega)^{n \times n}}, \end{aligned}$$

and then $||\nabla(u_{\varepsilon} - u_{PP})||_{L^{2}(\Omega)^{n}} \leq (c/\varepsilon)||\operatorname{div} u_{PP}||_{H^{-1}(\Omega)}$ follows.

Corollary 3.2. If u_{PP} satisfies div $u_{PP} = 0$, then $u_{\varepsilon} = u_{PP}$ and $p_{\varepsilon} = p_{PP}$ hold for all $\varepsilon > 0$. Furthermore, $u_S = u_{\varepsilon} = u_{PP}$ and $p_S = [p_{\varepsilon}] = [p_{PP}]$ hold for all $\varepsilon > 0$.

4 Links between (ES) and (S)

Lemma 4.1. If $v \in H^1(\Omega)^n$, $q \in L^2(\Omega)$ and $f \in H^{-1}(\Omega)^n$ satisfy

$$\int_{\Omega} \nabla v : \nabla \varphi + \langle \nabla q, \varphi \rangle = \langle f, \varphi \rangle \quad for \ all \ \varphi \in H^1_0(\Omega)^n,$$

then there exists a constant c > 0 such that

 $||q||_{L^{2}(\Omega)/\mathbb{R}} \leq c(||\nabla v||_{L^{2}(\Omega)^{n \times n}} + ||f||_{H^{-1}(\Omega)^{n}}).$

Proof. Let c be the constant arising in Theorem 2.2. Then we have

$$\begin{split} ||q||_{L^{2}(\Omega)/\mathbb{R}} &\leq c ||\nabla q||_{H^{-1}(\Omega)^{n}} = c \sup_{\varphi \in S_{n}} |\langle \nabla q, \varphi \rangle| \\ &\leq c \sup_{\varphi \in S_{n}} (|\int_{\Omega} \nabla v : \nabla \varphi| + |\langle f, \varphi \rangle|) \\ &\leq c (||\nabla v||_{L^{2}(\Omega)^{n \times n}} + ||f||_{H^{-1}(\Omega)^{n}}). \end{split}$$

Theorem 4.2. There exists a constant c > 0 independent of ε such that

$$||u_{\varepsilon}||_{H^{1}(\Omega)^{n}} \leq c, \quad ||p_{\varepsilon}||_{L^{2}(\Omega)/\mathbb{R}} \leq c \quad for \ all \ \varepsilon > 0.$$

Furthermore, we have

$$u_{\varepsilon} - u_{S} \rightharpoonup 0$$
 weakly in $H_{0}^{1}(\Omega)^{n}$, $[p_{\varepsilon}] - p_{S} \rightharpoonup 0$ weakly in $L^{2}(\Omega)/\mathbb{R}$ as $\varepsilon \rightarrow 0$.

Proof. We use the notations $u_0 \in H^1(\Omega)^n$, $p_0 \in H^1(\Omega)$, $f \in H^{-1}(\Omega)^n$ and $g \in H^{-1}(\Omega)$ in Theorem 2.6. We put $\tilde{u}_{\varepsilon} := u_{\varepsilon} - u_0 \in H^1_0(\Omega)^n$, $\tilde{p}_{\varepsilon} := p_{\varepsilon} - p_0 \in H^1_0(\Omega)$. Then we have

$$\begin{cases} \int_{\Omega} \nabla \tilde{u}_{\varepsilon} : \nabla \varphi + \int_{\Omega} (\nabla \tilde{p}_{\varepsilon}) \cdot \varphi = \langle f, \varphi \rangle & \text{for all } \varphi \in H_0^1(\Omega)^n, \\ \varepsilon \int_{\Omega} \nabla \tilde{p}_{\varepsilon} \cdot \nabla \psi + \int_{\Omega} (\operatorname{div} \tilde{u}_{\varepsilon}) \psi = \varepsilon \langle g, \psi \rangle & \text{for all } \psi \in H_0^1(\Omega). \end{cases}$$
(4.7)

Putting $\varphi := \tilde{u}_{\varepsilon}, \psi := \tilde{p}_{\varepsilon}$ and adding the two equations of (4.7), we get

$$||\nabla \tilde{u}_{\varepsilon}||^{2}_{L^{2}(\Omega)^{n\times n}} + \varepsilon||\nabla \tilde{p}_{\varepsilon}||^{2}_{L^{2}(\Omega)^{n}} \leq ||f||_{H^{-1}(\Omega)^{n}} ||\nabla \tilde{u}_{\varepsilon}||_{L^{2}(\Omega)^{n\times n}} + \varepsilon||g||_{H^{-1}(\Omega)} ||\nabla \tilde{p}_{\varepsilon}||_{L^{2}(\Omega)^{n}}$$

since $\int_{\Omega} (\nabla \tilde{p}_{\varepsilon}) \cdot \tilde{u}_{\varepsilon} = -\int_{\Omega} (\operatorname{div} \tilde{u}_{\varepsilon}) \tilde{p}_{\varepsilon}$. It leads that $(||\tilde{u}_{\varepsilon}||_{H^{1}(\Omega)^{n}})_{0 < \varepsilon < 1}$ and $(||\sqrt{\varepsilon} \tilde{p}_{\varepsilon}||_{H^{1}(\Omega)})_{0 < \varepsilon < 1}$ are bounded. In addition,

$$\|\tilde{p}_{\varepsilon}\|_{L^{2}(\Omega)/\mathbb{R}} \leq c(\|\nabla \tilde{u}_{\varepsilon}\|_{L^{2}(\Omega)^{n \times n}} + \|f\|_{H^{-1}(\Omega)^{n}})$$

by Lemma 4.1, which implies that $(||\tilde{p}_{\varepsilon}||_{L^{2}(\Omega)/\mathbb{R}})_{0<\varepsilon<1}$ is bounded. By Theorem 3.1, $(||u_{\varepsilon}||_{H^{1}(\Omega)^{n}})_{\varepsilon\geq1}$ and $(||\tilde{p}_{\varepsilon}||_{L^{2}(\Omega)/\mathbb{R}})_{\varepsilon\geq1}$ are bounded, and then $(||u_{\varepsilon}||_{H^{1}(\Omega)^{n}})_{\varepsilon>0}$ and $(||\tilde{p}_{\varepsilon}||_{L^{2}(\Omega)/\mathbb{R}})_{\varepsilon>0}$ are bounded.

Since $H_0^1(\Omega)^n \times (L^2(\Omega)/\mathbb{R})$ is reflexive and $(\tilde{u}_{\varepsilon}, [\tilde{p}_{\varepsilon}])_{0<\varepsilon<1}$ is bounded in $H_0^1(\Omega)^n \times (L^2(\Omega)/\mathbb{R})$, there exist $(u, p) \in H_0^1(\Omega)^n \times (L^2(\Omega)/\mathbb{R})$ and a subsequence of pair $(\tilde{u}_{\varepsilon_k}, \tilde{p}_{\varepsilon_k})_{k\in\mathbb{N}} \subset H_0^1(\Omega)^n \times H_0^1(\Omega)$ such that

$$\tilde{u}_{\varepsilon_k} \rightharpoonup u$$
 weakly in $H_0^1(\Omega)^n$, $[\tilde{p}_{\varepsilon_k}] \rightharpoonup p$ weakly in $L^2(\Omega)/\mathbb{R}$ as $k \to \infty$.

Hence, from (4.7) with $\varepsilon := \varepsilon_k$, taking $k \to \infty$, we obtain

$$\begin{cases} \int_{\Omega} \nabla u : \nabla \varphi + \langle \nabla p, \varphi \rangle = \langle f, \varphi \rangle & \text{for all } \varphi \in H_0^1(\Omega)^n \\ \int_{\Omega} (\operatorname{div} u) \psi = 0 & \text{for all } \psi \in H_0^1(\Omega), \end{cases}$$
(4.8)

where

$$\begin{split} |\varepsilon_k \int_{\Omega} \nabla \tilde{p}_{\varepsilon_k} \cdot \nabla \psi| &\leq \sqrt{\varepsilon_k} ||\sqrt{\varepsilon} \tilde{p}_{\varepsilon}||_{H^1(\Omega)} ||\psi||_{H^1(\Omega)} \to 0, \\ \int_{\Omega} \nabla \tilde{p}_{\varepsilon_k} \cdot \varphi &= -\int_{\Omega} [\tilde{p}_{\varepsilon_k}] \operatorname{div} \varphi \to -\int_{\Omega} p \operatorname{div} \varphi = \langle \nabla p, \varphi \rangle \end{split}$$

as $k \to \infty$. The first equation of (4.8) implies that

$$-\Delta(u+u_0) + \nabla(p+p_0) = F$$
 in $H^{-1}(\Omega)^n$.

From the second equation of (4.8), $\operatorname{div}(u+u_0) = 0$ follows. Hence, we obtain $u_S = u + u_0$ and $p_S = p + [p_0]$. Then we have

$$u_{\varepsilon_k} - u_S = u_{\varepsilon_k} - u - u_0 = \tilde{u}_{\varepsilon_k} - u_S \rightharpoonup 0 \text{ weakly in } H_0^1(\Omega)^n,$$

$$[p_{\varepsilon_k}] - p_S = [p_{\varepsilon_k} - p - p_0] = [\tilde{p}_{\varepsilon_k}] - p \rightharpoonup 0 \text{ weakly in } L^2(\Omega)/\mathbb{R}.$$

An arbitrarily chosen subsequence of $((u_{\varepsilon}, [p_{\varepsilon}]))_{0 < \varepsilon < 1}$ has a subsequence which converges to (u_S, p_S) , so we can conclude the proof.

Theorem 4.3. Suppose that $p_S \in H^1(\Omega)$. Then we have

$$u_{\varepsilon} - u_S \to 0 \text{ strongly in } H_0^1(\Omega)^n, \ [p_{\varepsilon}] - p_S \to 0 \text{ strongly in } L^2(\Omega)/\mathbb{R} \text{ as } \varepsilon \to 0$$

Proof. We have

$$\begin{cases} \int_{\Omega} \nabla(u_{\varepsilon} - u_{S}) : \nabla\varphi + \int_{\Omega} (\nabla(p_{\varepsilon} - p_{S})) \cdot \varphi = 0 \quad \text{for all } \varphi \in H_{0}^{1}(\Omega)^{n}, \\ \varepsilon \int_{\Omega} \nabla(p_{\varepsilon} - p_{S}) \cdot \nabla\psi + \int_{\Omega} (\operatorname{div} u_{\varepsilon})\psi = 0 \quad \text{for all } \psi \in H_{0}^{1}(\Omega). \end{cases}$$
(4.9)

We use the notations $p_0 \in H^1(\Omega)$ in Theorem 2.6. Putting $\varphi := u_{\varepsilon} - u_S \in H^1_0(\Omega)^n, \psi := p_{\varepsilon} - p_0 \in H^1_0(\Omega)$ and $\tilde{p}_S := p_S - p_0 \in H^1(\Omega)$, we get

$$\begin{aligned} ||\nabla(u_{\varepsilon} - u_{S})||^{2}_{L^{2}(\Omega)^{n \times n}} + \varepsilon ||\nabla(p_{\varepsilon} - p_{S})||^{2}_{L^{2}(\Omega)^{n}} \\ &= \int_{\Omega} (\nabla \tilde{p}_{S}) \cdot (u_{\varepsilon} - u_{S}) - \varepsilon \int_{\Omega} \nabla(p_{\varepsilon} - p_{S}) \cdot \nabla \tilde{p}_{S} \\ &\leq ||\nabla \tilde{p}_{S}||_{L^{2}(\Omega)^{n}} ||u_{\varepsilon} - u_{S}||_{L^{2}(\Omega)^{n}} + \varepsilon ||\nabla(p_{\varepsilon} - p_{S})||_{L^{2}(\Omega)^{n}} ||\nabla \tilde{p}_{S}||_{L^{2}(\Omega)^{n}}. \end{aligned}$$

By Corollary 4.2 and the Rellich-Kondrachov Theorem, there exists a sequence $(\varepsilon_k)_{k\in\mathbb{N}}\subset\mathbb{R}$ such that

$$u_{\varepsilon_k} \to u_S$$
 strongly in $L^2(\Omega)^n$.

So, we can write that

$$\begin{aligned} &||\nabla(u_{\varepsilon_k} - u_S)||^2_{L^2(\Omega)^{n \times n}} \\ \leq &||\nabla \tilde{p}_S||_{L^2(\Omega)^n} ||u_{\varepsilon_k} - u_S||_{L^2(\Omega)^n} + \sqrt{\varepsilon_k} ||\sqrt{\varepsilon_k} \nabla(p_{\varepsilon_k} - p_S)||_{L^2(\Omega)^n} ||\nabla \tilde{p}_S||_{L^2(\Omega)^n}. \end{aligned}$$

It implies that

$$||[p_{\varepsilon_k}] - p_S||_{L^2(\Omega)/\mathbb{R}} = ||p_{\varepsilon_k} - p_S||_{L^2(\Omega)/\mathbb{R}} \le c||\nabla(u_{\varepsilon_k} - u_S)||_{L^2(\Omega)^{n \times n}} \to 0$$

by Lemma 4.1. An arbitrarily chosen subsequence of $((u_{\varepsilon}, [p_{\varepsilon}]))_{0 < \varepsilon < 1}$ has a subsequence which converges to (u_S, p_S) , so we can conclude the proof.

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References

- A. A. Amsden and F. H. Harlow, A simplified MAC technique for incompressible fluid flow calculations, J. Comput. Phys. 6 (1970), 322–325.
- [2] H. Brezis, Functional Analysis, Sobolev Spaces and Partial Differential Equations, Springer, 2010.
- [3] R. K.-C Chan and R. L. Street, A computer study of finite-amplitude water waves, J. Comput. Phys. 6 (1970), 68–94.
- [4] A. J. Chorin, Numerical solution of the Navier-Stokes equations, Math. Comput. 22 (1968), no. 104, 745–762.
- [5] C. Conca, F. Murat, and O. Pironneau, The Stokes and Navier-Stokes equations with boundary conditions involving the pressure, Jpn. J. Math. 20 (1994), no. 2, 279–318.
- [6] C. Conca, C. Pares, O. Pironneau, and M. Thiriet, Navier-Stokes equations with imposed pressure and velocity fluxes, Int. Numer. Meth. Fl. 20 (1995), 267–287.
- [7] S. J. Cummins and M. Rudman, An SPH projection method, J. Comput. Phys. 152 (1999), 584–607.
- [8] J. Douglas, Jr. and J. Wang, An absolutely stabilized finite element method for the Stokes problem, Math. Comput. 52 (1989), no. 186, 495–508.
- [9] G. Duvaut and J. L. Lions, Inequalities in Mechanics and Physics, Springer-Verlag, 1976.
- [10] V. Girault and P.-A. Raviart, Finite Element Methods for Navier-Stokes Equations, Springer-Verlag, 1986.
- [11] R. Glowinski, *Finite Element Methods for Incompressible Viscous Flow*, Handbook of Numerical Analysis (P. G. Ciarlet and J. L. Lions, eds.), vol. 9, North-Holland, 2003, pp. 3–1176.
- [12] J.-L. Guermond and L. Quartapelle, On stability and convergence of projection methods based on pressure Poisson equation, Int. J. Numer. Meth. Fluids 26 (1998), 1039–1053.
- [13] F. H. Harlow and J. E. Welch, Numerical calculation of time-dependent viscous incompressible flow of fluid with a free surface, The Physics of Fluids 8 (1965), 2182–2189.

- [14] T. J. R. Hughes, L. P. Franca, and M. Balestra, A new finite element formulation for computational fluid dynamics: V. circumventing the Babuška-Brezzi condition: A stable Petrov-Galerkin formulation of the Stokes problem accommodating equalorder interpolations, Comput. Methods Appl. Mech. Engrg. 59 (1986), 85–99.
- [15] J. Kim and P. Moin, Application of a fractional-step method to incompressible Navier-Stokes equations, J. Comput. Phys. 59 (1985), 308–323.
- [16] S. Marušić, On the Navier-Stokes system with pressure boundary condition, Ann. Univ. Ferrara 53 (2007), 319–331.
- [17] S. McKee, M. F. Tomé, J. A. Cuminato, A. Castelo, and V. G. Ferreira, Recent advances in the marker and cell method, Arch. Comput. Meth. Engng. 2 (2004), 107–142.
- [18] J. Nečas, Les méthodes directes en théorie des équations elliptiques, Academia, Praha, and Masson et Cie, Editeurs, Paris, 1967.
- [19] J. B. Perot, An analysis of the fractional step method, J. Comput. Phys. 108 (1993), 51–58.
- [20] R. Temam, Navier-Stokes Equations, North Holland, 1979.
- [21] J. A. Viecelli, A computing method for incompressible flows bounded by moving walls, J. Comput. Phys. 8 (1971), 119–143.