# NUMERICAL EXPERIMENTS ON ANALYTICITY OF SOLUTIONS TO FRACTIONAL DIFFERENTIAL EQUATIONS 

Hiroshi Fujiwara<br>Graduate School of Informatics, Kyoto University, Yoshida-honmachi, Sakyo-ku, Kyoto, 606-8501, Japan<br>(fujiwara@acs.i.kyoto-u.ac.jp)<br>Nobuyuki Higashimori<br>Center for the Promotion of Interdisciplinary Education and Research, Kyoto University,<br>Yoshida-honmachi, Sakyo-ku, Kyoto, 606-8501, Japan<br>(nobuyuki@acs.i.kyoto-u.ac.jp)<br>and<br>Hitoshi Imai<br>Faculty of Science and Engineering, Doshisha University, 1-3 Tatara Miyakodani, Kyotanabe, Kyoto, 610-0394, Japan<br>(himai@mail.doshisha.ac.jp)


#### Abstract

In the paper, analyticity of solutions of fractional differential equations is investigated numerically. In numerical computation the Chebyshev spectral collocation method is used for discretization to realize arbitrary order approximation. Multipleprecision arithmetic is also used for arbitrary reduction of rounding error. Some onedimensional fractional differential equations are solved numerically. Numerical results clarify analyticity of solutions.


Communicated by Editors; Received April 17, 2018
This work was supported by JSPS KAKENHI Grant Numbers 16K13774 and 16H02155.
AMS Subject Classification: 65-05, 65L05.
Keywords: fractional differential equation, analyticity, spectral method, multiple-precision arithmetic.

## 1 Introduction

Recently, there has been a tremendous increase in the use of fractional differential models to simulate dynamics in many fields, e.g. physics, chemistry, biology, engineering and so on $[1],[4],[6],[7],[17]$. In the models fractional order derivatives are used. They have a long history and they have been studied in the field of fractional calculus. A list of prominent mathematicians contributing to fractional calculus is found in [20]. Besides them, Caputo gave a convenient definition for initial value problems of differential equations as follows [5]: For $0<\alpha$ non-integer,

$$
\begin{equation*}
{ }_{\mathrm{c}} D_{0}^{\alpha} u(t)=\frac{1}{\Gamma(n-\alpha)} \int_{0}^{t}(t-s)^{n-\alpha-1} \frac{d^{n} u(s)}{d s^{n}} d s \tag{1}
\end{equation*}
$$

where $\Gamma$ is the gamma function, $n=\lceil\alpha\rceil$ is the least integer more than or equal to $\alpha$, and $u$ is assumed to be sufficiently smooth. This definition is used in the paper because it is convenient for our numerical computation.

It is hard to obtain exact solutions of fractional differential equations, or even if an exact solution is available it may be too complicated to be used in practice. Thus the role of numerical computation becomes important. Various numerical methods to fractional differential equations such as the Adams-type predictor-corrector method, linear multistep methods, are presented in [1] and [7]. Spectral methods have also been applied to fractional differential equations [7],[9],[21]. As for spectral methods it is well known that they are superior in accuracy [3]. If a function $u$ is $m$-times continuously differentiable, then $\left\|u-u_{N}\right\|=O\left(N^{-m}\right)$ where $u_{N}$ is the $N$-th order approximate function by the spectral methods, and $\|\cdot\|$ is a suitable norm. This property is called spectral accuracy or infinite-order accuracy. Moreover, if a function $u$ is analytic in a strip of the complex plane, then $\left\|u-u_{N}\right\|=O\left(e^{-\gamma N}\right)$ with some positive number $\gamma$. In [21] numerical results for some fractional differential equations show spectral accuracy.

Spectral methods are roughly divided into the Galerkin type and the collocation type [2],[3]. The latter type is called the spectral collocation method, and it is more convenient for nonlinear problems or higher dimensional problems. In the paper the Chebyshev spectral collocation method is adopted. Although in [7] and [21] spectral collocation methods are already used, by using Chebyshev polynomials we can give a concrete expression of coefficients in the linear combination of basis functions. In the paper multiple-precision arithmetic is also used in numerical computation. To investigate mathematical properties of solutions numerically it is necessary to eliminate the influence of rounding error. The combination of spectral methods and multiple-precision arithmetic was presented in [16] to realize the arbitrary reduction of numerical error. Such reduction is indispensable for direct numerical simulations of inverse problems $[10],[13],[15]$. By using this combination we also investigated numerically mathematical properties of functions or solutions, for instance, regularity or existence of solutions [8],[14],[19]. In the paper, by using the combination we investigate analyticity of solutions of fractional differential equations.

## 2 The Chebyshev Spectral Collocation Method

We derive a convenient form of fractional derivatives in terms of the Chebyshev spectral method, which plays an essential role in our numerical studies.

First we briefly recall the Chebyshev spectral collocation method [2]. Let $\tilde{U}_{N}(x)$ be a polynomial function of degree $N$ defined on the interval $[-1,1]$. The Chebyshev expansion of $\tilde{U}_{N}(x)$ is

$$
\begin{equation*}
\tilde{U}_{N}(x)=\sum_{k=0}^{N} \bar{a}_{k} T_{k}(x), \tag{2}
\end{equation*}
$$

where $T_{k}(x)=\cos (k \arccos x)$ is the Chebyshev polynomial of order $k$. We also introduce Chebyshev-Gauss-Lobatto collocation points as follows:

$$
\begin{equation*}
\bar{x}_{j}=\cos \frac{j \pi}{N} \quad(j=0,1,2, \cdots, N) . \tag{3}
\end{equation*}
$$

Letting $\tilde{U}_{j}=\tilde{U}_{N}\left(\bar{x}_{j}\right)$, from (2) we have

$$
\tilde{U}_{j}=\sum_{k=0}^{N} \bar{a}_{k} T_{k}\left(\bar{x}_{j}\right) .
$$

The coefficients $\bar{a}_{k}$ are given by the following inverse formula

$$
\begin{equation*}
\bar{a}_{k}=\frac{2}{N \bar{c}_{k}} \sum_{j=0}^{N} \frac{1}{\bar{c}_{j}} T_{k}\left(\bar{x}_{j}\right) \tilde{U}_{j}, \quad(k=0,1,2, \ldots, N) \tag{4}
\end{equation*}
$$

where

$$
\bar{c}_{j}= \begin{cases}2, & j=0, N \\ 1, & \text { otherwise }\end{cases}
$$

Substituting (4) to $\bar{a}_{k}$ in (2) we obtain

$$
\tilde{U}_{N}(x)=\sum_{j=0}^{N}\left(\frac{1}{\bar{c}_{j}} \sum_{k=0}^{N} \frac{2}{N \bar{c}_{k}} T_{k}\left(\bar{x}_{j}\right) T_{k}(x)\right) \tilde{U}_{j} .
$$

This representation motivates an $N$-th order approximation $U_{N}(x)$ to a function $U(x)$ on $[-1,1]$ as

$$
\begin{equation*}
U_{N}(x)=\sum_{j=0}^{N}\left(\frac{1}{\bar{c}_{j}} \sum_{k=0}^{N} \frac{2}{N \bar{c}_{k}} T_{k}\left(\bar{x}_{j}\right) T_{k}(x)\right) U_{j}, \tag{5}
\end{equation*}
$$

where $U_{j}=U\left(\bar{x}_{j}\right)$. We should remark that if $U(x)$ is a polynomial of degree $m(\leq N)$ then $U(x)=U_{N}(x)$. The truncation number $N$ in (5) denotes the approximation order. On the other hand, $N$ in (3) denotes the number of collocation points; to be exact, the number is $(N+1)$. This means that in spectral collocation methods the approximation order can be easily controlled by the number of collocation points. As stated in Introduction,
spectral collocation methods are applicable to nonlinear problems or higher dimensional problems. These features are quite convenient to investigate mathematical properties of solutions to various problems.

In spectral collocation methods, differential equations are satisfied at the collocation points. Thus they require the derivative of the approximate function at each collocation point. For instance, the first order derivative of $U_{N}(x)$ at the collocation point $\bar{x}_{i}$ is given by differentiating (5):

$$
\begin{equation*}
\frac{d}{d x} U_{N}\left(\bar{x}_{i}\right)=\sum_{j=0}^{N}\left(D_{x}\right)_{i, j} U_{j}=\sum_{j=0}^{N}\left(\frac{1}{\bar{c}_{j}} \sum_{k=0}^{N} \frac{2}{N \bar{c}_{k}} T_{k}\left(\bar{x}_{j}\right) \frac{d}{d x} T_{k}\left(\bar{x}_{i}\right)\right) U_{j} \tag{6}
\end{equation*}
$$

where $D_{x}$ is called the first order derivative matrix.
Let $u(t)$ be a function defined on $[0, T]$. By substituting $t=\frac{T}{2}(x+1), u\left(\frac{T}{2}(x+1)\right)$ is a function on $[-1,1]$ and is approximated by the Chebyshev expansion

$$
U_{N}(x)=\left(\begin{array}{lll}
T_{0}(x) & \cdots & T_{N}(x)
\end{array}\right)\left(\begin{array}{c}
\bar{a}_{0} \\
\vdots \\
\bar{a}_{N}
\end{array}\right) .
$$

By using a square matrix $\tilde{T}=\left(\tilde{T}_{k j}\right)$ of order $N+1$ with entries $\tilde{T}_{k j}=\frac{2}{N \bar{c}_{k} \bar{c}_{j}} T_{k}\left(\bar{x}_{j}\right)$, the inverse formula (4) in the matrix form is

$$
\left(\begin{array}{c}
\bar{a}_{0} \\
\vdots \\
\bar{a}_{N}
\end{array}\right)=\tilde{T}\left(\begin{array}{c}
U_{0} \\
\vdots \\
U_{N}
\end{array}\right)
$$

where $U_{j}=u\left(\frac{T}{2}\left(\bar{x}_{j}+1\right)\right)$. Let $C$ be a square matrix of order $N+1$ which satisfies

$$
C\left(\begin{array}{c}
1  \tag{7}\\
\frac{T}{2}(x+1) \\
\vdots \\
\left(\frac{T}{2}\right)^{N}(x+1)^{N}
\end{array}\right)=\left(\begin{array}{c}
T_{0}(x) \\
T_{1}(x) \\
\vdots \\
T_{N}(x)
\end{array}\right)
$$

We remark that $\tilde{T}$ and $C$ are constant matrices. They lead to the expression

$$
U_{N}(x)=\left(\begin{array}{llll}
1 & \frac{T}{2}(x+1) & \cdots & \left(\frac{T}{2}(x+1)\right)^{N}
\end{array}\right) C^{T} \tilde{T}\left(\begin{array}{c}
U_{0} \\
\vdots \\
U_{N}
\end{array}\right)
$$

where $C^{T}$ is the transpose of $C$. Let $t_{j}=\frac{T}{2}\left(\bar{x}_{j}+1\right), u_{j}=u\left(t_{j}\right)$, and $u_{N}(t)=U_{N}(x)$ with $t=\frac{T}{2}(x+1)$. Then $u_{j}=U_{j}$ and the function $u_{N}$ is an approximation to $u$ and is written
as

$$
u_{N}(t)=\left(\begin{array}{llll}
1 & t & \cdots & t^{N}
\end{array}\right) C^{T} \tilde{T}\left(\begin{array}{c}
u_{0} \\
\vdots \\
u_{N}
\end{array}\right) .
$$

This immediately leads to the fractional derivative of $u_{N}$ at $t_{i}$ as

$$
{ }_{\mathrm{c}} D_{0}^{\alpha} u_{N}\left(t_{i}\right)=\left(\begin{array}{llll}
{ }_{\mathrm{c}} D_{0}^{\alpha}[1]\left(t_{i}\right) & { }_{\mathrm{c}} D_{0}^{\alpha}[t]\left(t_{i}\right) & \cdots & \left.{ }_{\mathrm{c}} D_{0}^{\alpha}\left[t^{N}\right]\left(t_{i}\right)\right)
\end{array}\right) C^{T} \tilde{T}\left(\begin{array}{c}
u_{0} \\
\vdots \\
u_{N}
\end{array}\right),
$$

which is a linear combination of $u_{j}, j=0, \ldots, N$, and is similar to (6). It should be emphasized that the Caputo derivatives of the unknown function at each collocation points ${ }_{\mathrm{c}} D_{0}^{\alpha}\left[t^{k}\right]\left(t_{i}\right)$ are easily calculated since the Caputo derivatives of monomials $t^{k}$ for a nonnegative integer $k$ are given by

$$
{ }_{\mathrm{c}} D_{0}^{\alpha}\left[t^{k}\right]= \begin{cases}0, & 0 \leq k<\alpha \\ \frac{\Gamma(k+1)}{\Gamma(k+1-\alpha)} t^{k-\alpha}, & \alpha \leq k .\end{cases}
$$

It is noted in [21] that the entries of $C$ are calculated by the use of inverse of the Vandermonde matrix which is obtained by substituting $N+1$ collocation points $\bar{x}_{j}, j=$ $0,1, \ldots, N$, into (7). However, the entries of $C$ are calculated explicitly as follows. For the purpose, we introduce two square matrices $F=\left(f_{k j}\right)$ and $L=\left(\ell_{k j}\right)$ of order $N+1$, which satisfy

$$
\left(\begin{array}{c}
T_{0}(x) \\
\vdots \\
T_{N}(x)
\end{array}\right)=F\left(\begin{array}{c}
1 \\
x \\
\vdots \\
x^{N}
\end{array}\right) \quad \text { and } \quad\left(\begin{array}{c}
1 \\
x \\
\vdots \\
x^{N}
\end{array}\right)=L\left(\begin{array}{c}
1 \\
\frac{T}{2}(x+1) \\
\vdots \\
\left(\frac{T}{2}\right)^{N}(x+1)^{N}
\end{array}\right) .
$$

Firstly, the entries $f_{k j}$ are calculated from the recurrence relation $T_{0}(x)=1, T_{1}(x)=x$ and $T_{k+1}(x)=2 x T_{k}(x)-T_{k-1}(x)$. Secondly, it follows from the binomial expansion that

$$
x^{k}=(x+1-1)^{k}=\sum_{j=0}^{k}(-1)^{k-j}\binom{k}{j}(x+1)^{j} .
$$

Since $(x+1)^{j}, j=0, \ldots, N$, are linearly independent, $\ell_{k j}$ is given by

$$
\ell_{k j}= \begin{cases}(-1)^{k-j}\left(\frac{2}{T}\right)^{j}\binom{k}{j}, & j \leq k, \\ 0, & k<j\end{cases}
$$

Finally the product of two lower triangular matrices gives $C=F L$.

## 3 Numerical Evaluation of Fractional Derivatives

This section is devoted to investigate accuracy of numerical integration appearing in (1) and computational environments.

The integrand in (1) has a singularity at boundary $s=t$, and thus special attentions should be paid for the sake of reliable and accurate computation. The use of double exponential rule [18] is known as a novel method to overcome the difficulty. It is also convenient due to returning numerical results in given tolerance automatically.

We demonstrate an example with $0<\alpha<1$ and $u(t)=t^{2}$. Using the standard double exponential transformation $\phi(y)=\tanh \left(\frac{\pi}{2} \sinh y\right)$, the evaluation of integral at $t=1$ is

$$
{ }_{\mathrm{c}} D_{0}^{\alpha} u(1)=\frac{1}{\Gamma(1-\alpha)} \int_{-\infty}^{\infty} \frac{\phi(y)+1}{\left(\frac{1-\phi(y)}{2}\right)^{\alpha}} \frac{\phi^{\prime}(y)}{2} d y
$$

by introducing the discretization parameter $h>0$ and $y_{k}=k h, k \in \mathbb{Z}$, "the composite trapezoidal rule" gives an approximate series as

$$
\approx \frac{1}{\Gamma(1-\alpha)} \sum_{k \in \mathbb{Z}} \frac{\phi\left(y_{k}\right)+1}{\left(\frac{1-\phi\left(y_{k}\right)}{2}\right)^{\alpha}} \frac{\phi^{\prime}\left(y_{k}\right)}{2} h
$$

which is truncated to a finite sum as

$$
\approx \frac{1}{\Gamma(1-\alpha)} \sum_{K_{0} \leq k \leq K_{1}} \frac{\phi\left(y_{k}\right)+1}{\left(\frac{1-\phi\left(y_{k}\right)}{2}\right)^{\alpha}} \frac{\phi^{\prime}\left(y_{k}\right)}{2} h
$$

with some integers $K_{0}$ and $K_{1}$.
As $|k| \rightarrow \infty, \phi^{\prime}\left(y_{k}\right)$ converges to the zero more rapidly than growth of the integrand caused by the singularity, and their products and accumulations are negligible for sufficiently large $|k|$. Therefore choice of $h, K_{0}$, and $K_{1}$ is crucial for accuracy of numerical integration. Table 1 and Table 2 show numerical evaluation of

$$
I_{h}=\frac{1}{\Gamma(1-\alpha)} \sum_{K_{0} \leq k \leq K_{1}} \frac{\phi\left(y_{k}\right)+1}{\left(\frac{1-\phi\left(y_{k}\right)}{2}\right)^{\alpha}} \frac{\phi^{\prime}\left(y_{k}\right)}{2} h
$$

with $\alpha=0.5,0.9$ respectively. In the tables, we show discretization and truncation parameters as $L=K_{0} h, U=K_{1} h$, and $K=K_{1}-K_{0}$. Values of the exact solution

$$
\left.{ }_{\mathrm{c}} D_{0}^{\alpha}\left[t^{2}\right]\right|_{t=1}=\frac{2}{\Gamma(3-\alpha)}
$$

are also shown in tables.
From Table 1, both double precision arithmetic and 100 decimal digit computation enable us to find approximate values with $(L, U)=(-5,3.16)$ for $\alpha=0.5$. On the other hand, for $\alpha=0.9$, parameters $(L, U)=(-5,5)$ and $K \geq 16$ in 100 decimal digit return approximations while the double precision arithmetic returns infinity if $U$ is greater

Table 1: Numerical integration $I_{h}$ with $\alpha=0.5$.

| $L$ | $U$ | $K$ | double precision | 100 digits |
| :--- | :--- | ---: | :---: | :---: |
| -5.00 | 3.16 | 8 | 1.50652114572 | 1.50652118788 |
|  |  | 16 | 1.50451382597 | 1.50451384705 |
|  |  | 32 | 1.50450554499 | 1.50450555553 |
| -5.00 | 3.17 | 8 | $+\infty$ | 1.50798169193 |
|  |  | 16 | $+\infty$ | 1.50451519723 |
|  |  | 32 | $+\infty$ | 1.50450555566 |
| exact value |  | $1.50450555613 \ldots$ |  |  |

Table 2: Numerical integration $I_{h}$ with $\alpha=0.9$.

| $L$ | $U$ | $K$ | double precision | 100 digits |
| :--- | :--- | ---: | :---: | :---: |
| -5.00 | 3.16 | 8 | 1.87072828925 | 1.90636895125 |
|  |  | 16 | 1.87900443987 | 1.89682476417 |
|  |  | 32 | 1.87188368373 | 1.88079405669 |
| -5.00 | 3.17 | 8 | $+\infty$ | 1.90645242453 |
|  |  | 16 | $+\infty$ | 1.89759689138 |
|  |  | 32 | $+\infty$ | 1.88210055004 |
| -5.00 | 5.00 | 8 | $+\infty$ | 1.90117136452 |
|  |  | 16 | $+\infty$ | 1.91113179851 |
|  |  | 32 | $+\infty$ | 1.91115819291 |
| exact value |  | $1.91115819293 \ldots$ |  |  |

than or equal to 3.17. This means that double precision arithmetic is not enough for numerical evaluation of fractional derivatives without special treatments of singularities in $I_{h}$, especially when the power $-\alpha$ is close to -1 . In our numerical computation, exflib [11] is used for multiple-precision arithmetic.

Figure 1 shows profiles of fractional derivatives ${ }_{\mathrm{c}} D_{0}^{\alpha} u(t)$ for $u(t)=t^{2}$ obtained by numerically integrating (1). Figure 1(a) and (b) are processed with double precision arithmetic and 100 decimal digits arithmetic respectively. We set $(L, U)=(-5,3.16)$ in double precision arithmetic, and $(L, U)=(-5,5)$ in 100 decimal digits arithmetic, with sufficiently large $K$. Obviously, results for $\alpha=0.9$ do not coincide. This also shows that double precision is not enough for reliable computations. More detail discussions and other numerical examples will be found in [12].


Figure 1: Fractional derivatives ${ }_{\mathrm{c}} D_{0}^{\alpha} u(t)$ of $u(t)=t^{2}$.

## 4 Numerical Results for Fractional Differential Equations

In this section we shall demonstrate some examples which shows efficiency of combination of the Chebyshev spectral method and multiple-precision arithmetic to achieve arbitrary accurate numerical solutions to fractional order differential equations.

## Example 1.

$$
\begin{aligned}
& u^{(1 / 3)}(t)=f(t), \quad 0<t<3, \\
& u(0)=0 .
\end{aligned}
$$

Three cases are considered: $f(t)=\frac{5!}{\Gamma(17 / 3)} t^{14 / 3}, \frac{10!}{\Gamma(32 / 3)} t^{29 / 3}$, and $\frac{15!}{\Gamma(47 / 3)} t^{44 / 3}$. The exact solutions are $u(t)=t^{5}, t^{10}$, and $t^{15}$ respectively. Since each solution is a monomial, it is exactly treated by the Chebyshev spectral method, and hence only rounding errors appear as numerical errors.

In Figure 2, the horizontal axis is the spectral degree $N$ and the vertical axis shows errors in numerical results:

$$
\max _{0 \leq j \leq N}\left|u\left(\bar{x}_{j}\right)-u_{N}\left(\bar{x}_{j}\right)\right|
$$

in the logarithmic scale. In the figure, purple, green, and blue graph represent numerical errors for $u(t)=t^{5}, t^{10}$, and $t^{15}$ respectively. In each color, plus $(+) \operatorname{sign}$ and cross $(\times)$ are error in 50 decimal digits and 100 decimal digits respectively. In each setting, behaviour of errors suddenly changes at a certain order $N_{0}$. More precisely, when $N$ becomes larger than or equal to $N_{0}$, the numerical errors are less than $10^{-50}$ or $10^{-100}$ which correspond to the rounding errors in each computation. This indicates that the solution is a polynomial of degree almost $N_{0}$.


Figure 2: Error decay for Example 1, the exact solutions are monomials.

## Example 2. [9]

$$
\begin{aligned}
& u^{(5 / 2)}(t)-3 u^{(2 / 3)}(t)=f(t), \quad 0<t \leq 1, \\
& u(0)=1, \quad u^{\prime}(0)=\beta, \quad u^{\prime \prime}(0)=\beta^{2} .
\end{aligned}
$$

We set the exact solution $u(t)=e^{\beta t}$, and calculate $f(t)$ as the left-hand side of the equation by using (1). Figure 3 shows numerical results by double precision (blue), 50 decimal digits (green), and 100 decimal digits (purple). The error decreases exponentially until the level of rounding error is almost reached. This behavior of errors is expected from the property of the spectral methods.


Figure 3: Error decay for Example 2, the exact solution is $e^{\beta t}$.

## Example 3. [9]

$$
\begin{aligned}
& u^{\prime \prime}(t)+u^{(1.5)}(t)+u(t)=f(t), \quad 0<t \leq 1, \\
& u(0)=0, \quad u^{\prime}(0)=\omega .
\end{aligned}
$$

The equation is known as the Bagley-Torvik equation. We set the exact solution $u(t)=$ $\sin \omega t$. Then $f(t)$ is calculated by using (1). Numerical results in several precisions are shown in Figure 4. The error decreases similarly in the preceding example, since the exact solution is also analytic in this example.


Figure 4: Error decay for Example 3, the exact solution is $\sin \omega t$.

Example 4. In the same equation and initial condition as Example 1, we adopt $f(t)=$ $\Gamma(7 / 3) t$, and the exact solution is $u(t)=t^{4 / 3}$. Figure 5 shows numerical results, where both the horizontal and vertical axes are in the logarithmic scale. In this setting numerical errors decay as $O\left(N^{-2.66}\right)$, which is not exponential decay but polynomial decay with respect to $N$. This suggests that the exact solution is not analytic.

## 5 Conclusion

In the paper, by using the Chebyshev spectral collocation method a concrete expression of coefficients in the linear combination of basis functions is derived. It is used for solving one-dimensional fractional differential equations numerically. We also investigate analyticity of solutions numerically. To do so, in numerical computation the multiple-precision arithmetic is used to estimate the influence of rounding error. Numerical results are very satisfactory. Spectral accuracy is seen and it shows analyticity of solutions. Moreover, when exact solutions are monomials, numerical results are succeeded to determine the degree of monomials.


Figure 5: Error decay for Example 4 in double precision arithmetic (Blue) and 100 decimal digits (Green), the exact solution is $t^{4 / 3}$.

## References

[1] D. Baleanu et al., Fractional Calculus: Models and Numerical Methods, World Scientific, Singapore, 2012.
[2] C. Canuto et al., Spectral Methods: Fundamentals in Single Domains, Springer, Berlin, 2006.
[3] C. Canuto et al., Spectral Methods: Evolution to Complex Geometries and Applications to Fluid Dynamics, Springer, Berlin, 2007.
[4] R. Caponetto et al., Fractional Order Systems: Modeling and Control Applications, World Scientific, Singapore, 2010.
[5] M. Caputo, Linear Model of Dissipation Whose $Q$ is Almost Frequency Independent - II, Geophys J. R. Astron. Soc., 13(1967), pp. 529-539.
[6] C. Cattani et al., Fractional Dynamics, De Gruyter Open, Berlin, 2015.
[7] A. Dabiri and E. A. Butcher, Efficient Modified Chebyshev Differentiation Matrices for Fractional Differential Equations, Commun. Nonlinear Sci. Numer. Simulat., 50(2017), pp. 284-310.
[8] K. C. Datta, H. Imai and H. Sakaguchi, On numerical Challenge for the Distinction between Smooth Functions and Analytic Functions, Theo. Appl. Mech. Jpn., 62(2013), pp. 107-117.
[9] E. H. Doha, A. H. Bhrawy and S. S. Ezz-Eldien, Efficient Chebyshev Spectral Methods for Solving Multi-term Fractional Orders Differential Equations, Appl. Math. Model., 35(2011), pp. 5662-5672.
[10] H. Fujiwara et al., High-precision Numerical Computation of Integral Equations of the First Kind, Transactions of JSIAM, 15(3) (2005), pp. 419-434 (in Japanese).
[11] H. Fujiwara and Y. Iso, Design of a Multiple-precision Arithmetic Package for a 64-bit Computing Environment and Its Application to Numerical Computation of Ill-posed Problems, IPSJ J., 44(3) (2003), pp. 925-931.
[12] H. Fujiwara, Multiple-precision Arithmetic Environment in MATLAB and Its Application to Reliable Computation of Fractional Order Derivatives, Math for Industry Kenkyu, 8 (2018) pp. 83-120
[13] H. Imai et al., A Direct Approach to an Inverse Problem, GAKUTO Int. Ser. Math. Sci. Appl., 12 (1999), pp. 223-232.
[14] H. Imai, S. Sakaguchi and Y. Iso, On Numerical Computation of the Toricomi Equation, Theo. Appl. Mech. Jpn., 59 (2010), pp. 359-372.
[15] H. Imai and T. Takeuchi, Some Advanced Applications of the Spectral Collocation Method, GAKUTO Int. Ser. Math. Sci. Appl., 17 (2001), pp. 323-335.
[16] H. Imai, T. Takeuchi and M. Kushida, On Numerical Simulation of Partial Differential Equations in Infinite Precision, Adv. Math. Sci. Appl., 9(2) (1999), pp. 10071016.
[17] R. L. Magin, Fractional Calculus in Bioengineering, Begell House, Connecticut, 2006.
[18] H. Takahashi and M. Mori, Double Exponential Formulas for Numerical Integration, Publ. Res. Inst. Math. Sci., Kyoto Univ., 9 (1974), pp. 721-741.
[19] T. Takeuchi and H. Imai, Direct Numerical Simulations of Cauchy Problems for the Laplace Operator, Adv. Math. Sci. Appl., 13(2) (2003), pp. 587-609.
[20] J. A. Tenreiro Machado, V. Kiryakova, and F. Mainardi, A Poster about the Recent History of Fractional Calculus, Fract. Calc. Appl. Anal., 13(3)(2010), pp. 329-334.
[21] W. Y. Tian, W. Deng and Y. Wu, Polynomial Spectral Collocation Method for Space Fractional Advection Diffusion Equation, Numer. Methods Partial Differ. Equ., 30(2)(2014), pp. 514-535.

